

# Meson Production in Antinucleon Annihilation on Nuclei

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Deutsche  
Forschungsgemeinschaft  
**DFG**

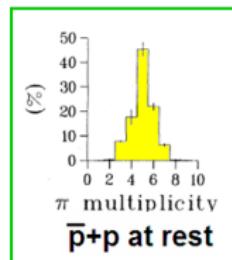
## upcoming experiments at the Facility of Antiproton Ion Research

- Anti-Proton **ANnihilation** at **DArmstadt**
- Antiproton-Ion-Collider



## what has been studied:

- low energy antiproton-nucleus scattering  
(Bachelor thesis: Thorsten Steinert)
- energy spectra of antiprotonic atoms  
(Bachelor thesis: Jan Haas)
- meson multiplicities



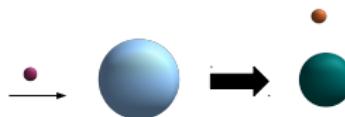


$$d^9\sigma_{\alpha\beta} = N_{\alpha\beta} \left(\frac{\hbar c}{2\pi}\right)^9 \frac{d^3k_1}{E_1} \frac{d^3k_2}{E_2} \frac{d^3k_B}{E_B} \left| M_{\alpha\beta} (\vec{k}_1, \vec{k}_2, \vec{k}_B; \vec{k}_\alpha) \right|^2 \\ \delta(\vec{k}_1 + \vec{k}_2 + \vec{k}_B) \delta(E_1 + E_2 + E_B - \sqrt{s})$$

$$M_{\alpha\beta} \approx t_{\bar{N}N \rightarrow 2\pi}(s) \langle \chi_{1\beta}^{(-)} \chi_{2\beta}^{(-)} | \varphi_B | \chi_{\bar{N}A}^{(+)} \rangle$$

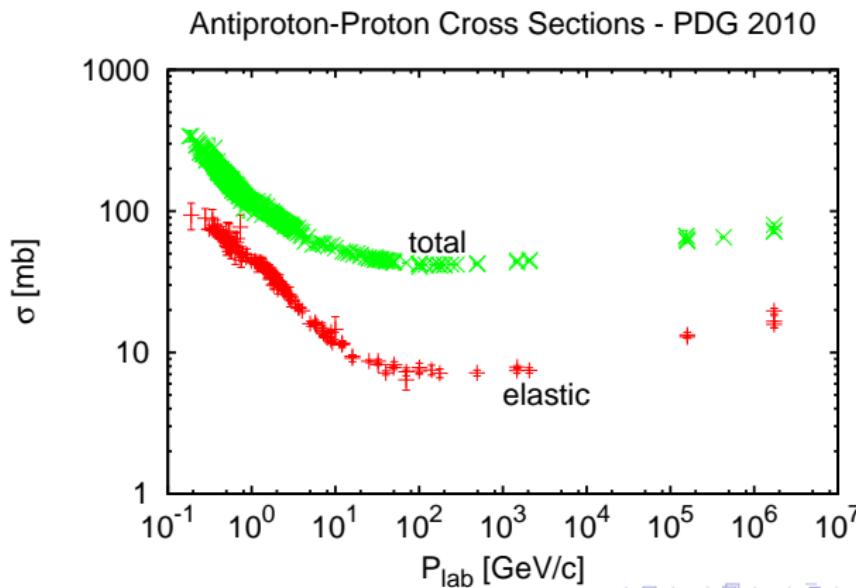
(1)

$$\varphi_B = \langle B | \psi_N | A \rangle = \sum_i \varphi_i \langle B | a_i | A \rangle$$
(2)



# PDG Cross Section

$$U_{opt}(\mathbf{r}) = V - iW = \sum_{N=p,n} \int \frac{d^3q}{(2\pi)^3} \rho_N(q) t_{\bar{p}N}(T_{Lab}, q^2) e^{i\mathbf{q}\cdot\mathbf{r}}$$



**elastic part:**

G-parity-transformation of the  $NN$  (Paris [1], Bonn [2]) interaction (charge conjugation plus  $180^\circ$  rotation around the y axis in isospin space):

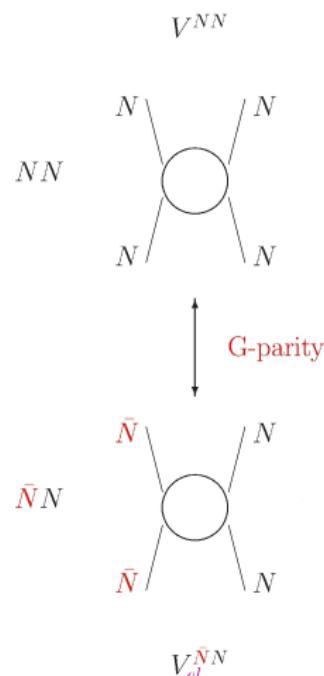
**odd G-Parity vertices**

$$V_{NN}(\pi, \omega, \delta) = -V_{\bar{N}N}(\pi, \omega, \delta)$$

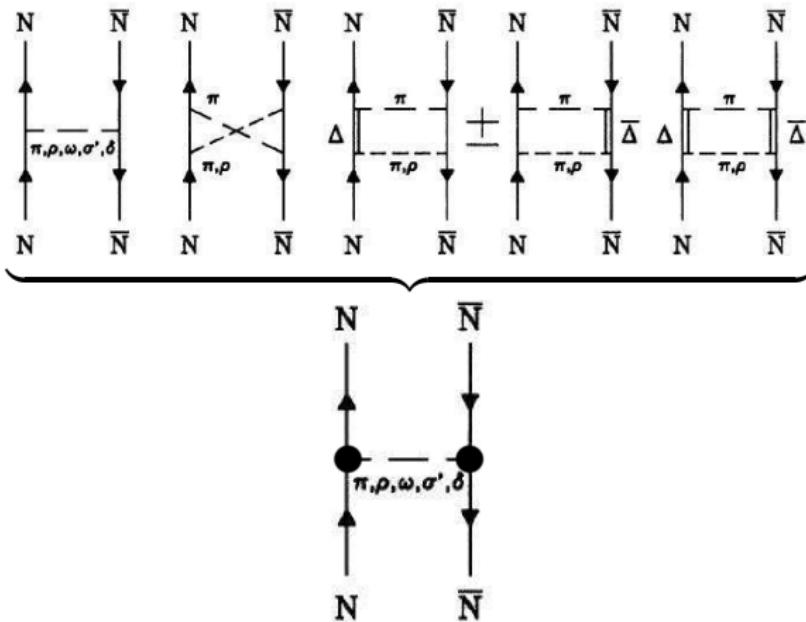
**even G-parity vertices**

$$V_{NN}(\sigma, \rho, \eta) = V_{\bar{N}N}(\sigma, \rho, \eta)$$

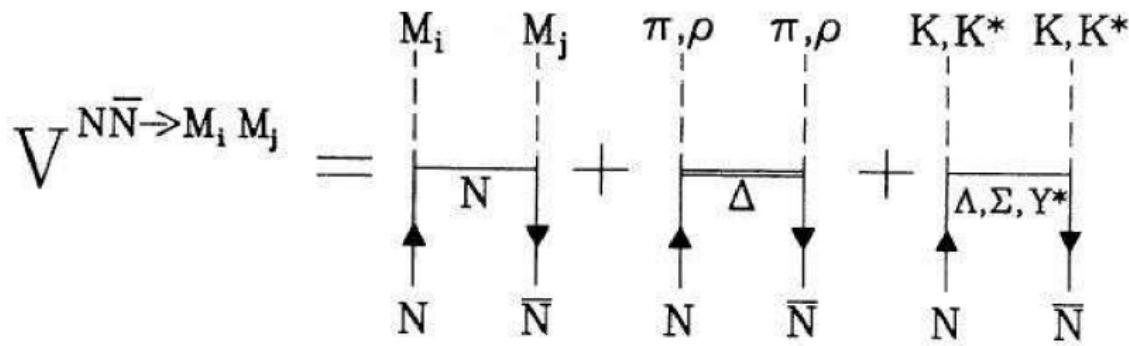
[1] Phys. Rev C 79 (2009) 054001, [2] Phys. Rev. C 51 (1995) 2360



# 2-Meson Diagrams Included in Bonn Model (elastic)

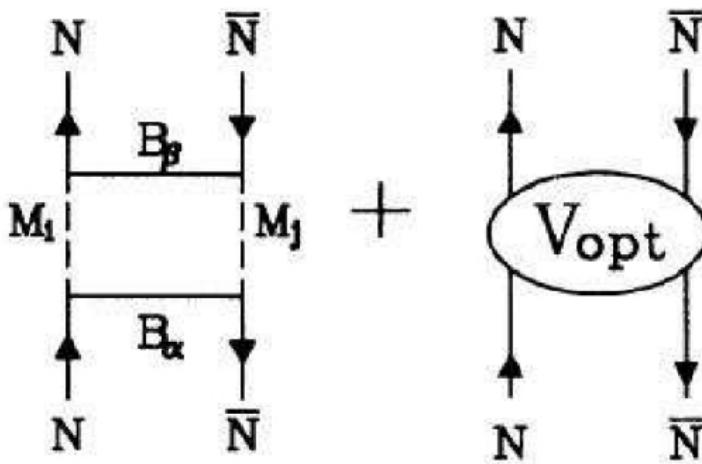


# Diagrams Included in Bonn Model (annihilation)



$$M_{i,j} = \pi, \eta, \rho, \omega, f_0, a_0, f_1, a_1, f_2, a_2$$

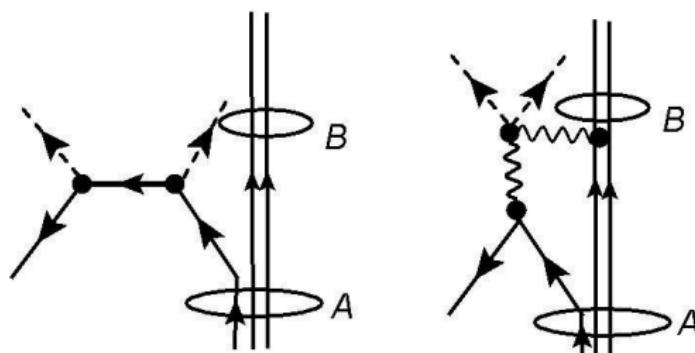
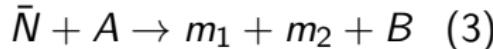
# Diagrams Included in Bonn Model (dispersive)



## 2-Meson Production on a Nucleus

### 2-Meson Production Recipe

- $\bar{N}A$  interaction
- Production Vertex
- $Bm_1m_2$  interaction



# Final State Interaction

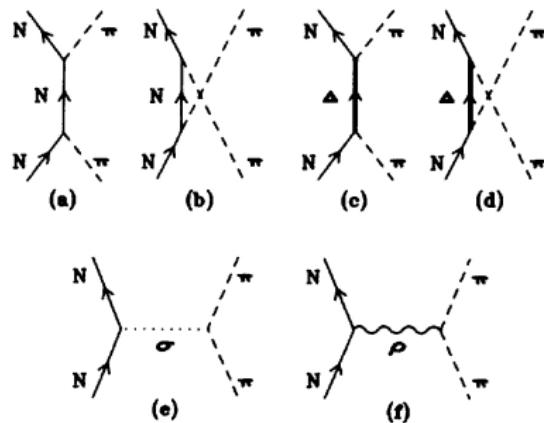


FIG. 1. Direct (a,c) and exchange (b,d) nucleon ( $N$ ) and delta-isobar ( $\Delta$ ) pole diagrams together with  $\sigma, \rho$ -exchange processes (e,f) used in  $\pi N$  interaction models.

# Final State Interaction

Pion-nucleus potential of Kisslinger type:

$$[-\Delta - k^2 + U_s + \vec{\partial} \cdot U_p \cdot \vec{\partial}] \Phi = 0 \quad (4)$$

After Krell-Ericson transformation  $\Phi = (1 - U_p)^{-1/2} \psi$  the local potential is transformed into (Johnson and Satchler):

$$U_N(r) = \frac{(\hbar c)^2}{2\omega} \left\{ \frac{U_s}{1 - U_p} - \frac{k^2 U_p}{1 - U_p} - \left[ \frac{\frac{1}{2} \vec{\nabla}^2 U_p}{1 - U_p} + \left( \frac{\frac{1}{2} \vec{\nabla} U_p}{1 - U_p} \right)^2 \right] \right\} \quad (5)$$

Pion-Nucleus interactions beyond the  $\Delta$ -resonance require higher resonances.

relation between the optical potential and the scattering amplitude:

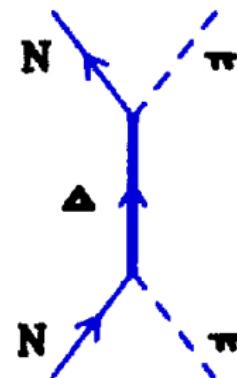
$$U(r) = 4\pi f_{\pi N} \rho(r) \quad (6)$$

$$U(r) = \underbrace{4\pi f_S \rho(r)}_{U_S} + \underbrace{4\pi 2f^P \frac{\partial}{\partial r} \rho(r) \frac{\partial}{\partial r}}_{U_P} \quad (7)$$

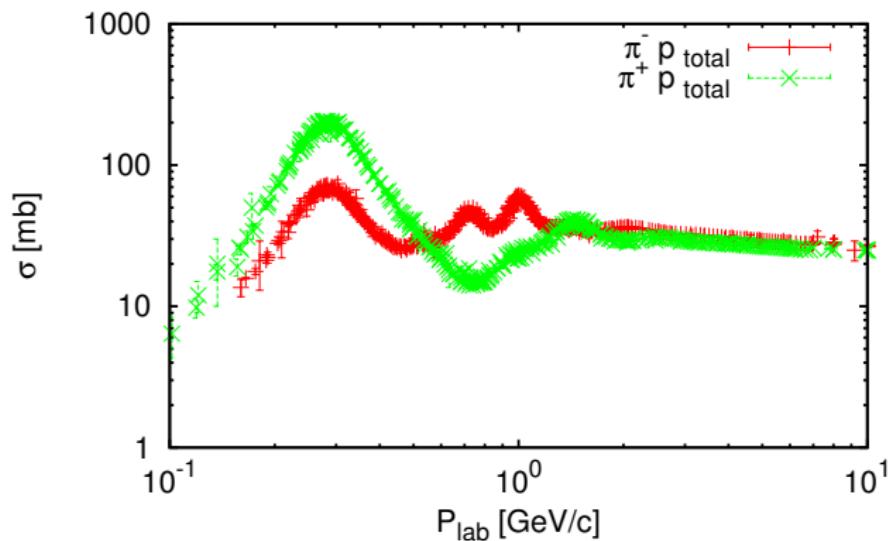
$$f^P(k, k') = \frac{\gamma v(k)v(k')}{(E_r - E - i\gamma k^3 [v(k)]^2)} \quad (8)$$

where the form-factor is

$$v(k) = \left[1 + \frac{k^2}{\kappa^2}\right]^{-2}, \quad (9)$$



## Pion-Proton Cross Sections - PDG 2010

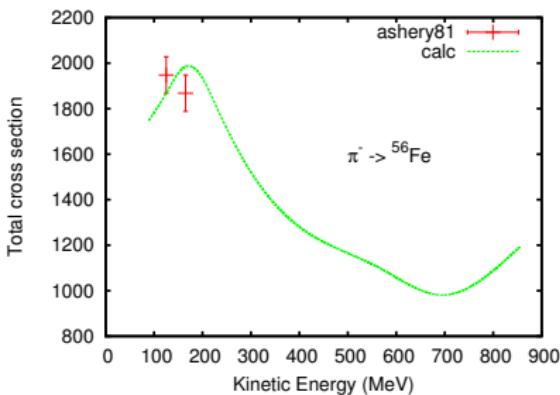
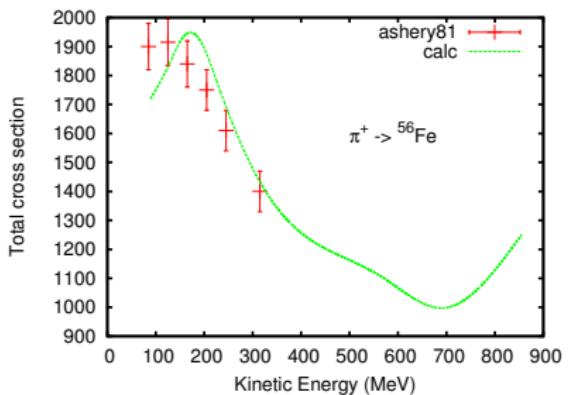
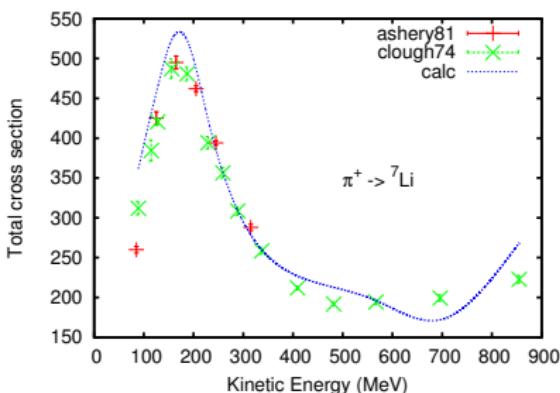
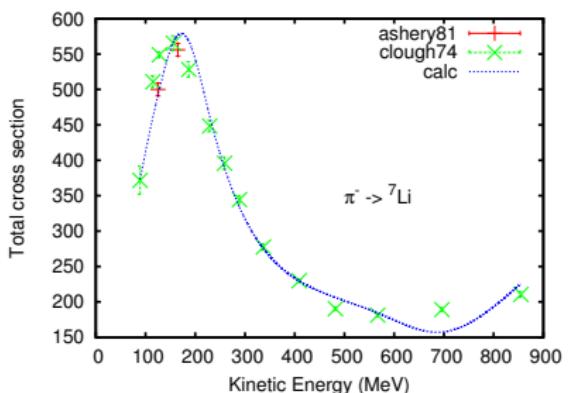


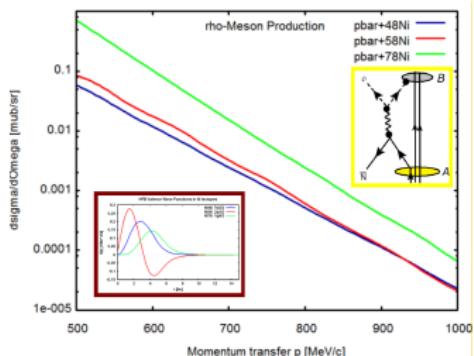
**Table:** Higher  $I = 1/2$  resonance parameters, PDF data.

| Name ( $E_r$ ) | $\Gamma$ [MeV] | R    | $I \ J^P$     | rank |
|----------------|----------------|------|---------------|------|
| $N(1440)$      | 300            | 0.6  | $1/2 \ 1/2^+$ | **** |
| $N(1535)$      | 150            | 0.35 | $1/2 \ 1/2^-$ |      |
| $N(1650)$      | 165            | 0.60 | $1/2 \ 1/2^-$ |      |
| $N(1675)$      | 150            | 0.60 | $1/2 \ 5/2^-$ | **** |
| $N(1680)$      | 130            | 0.68 | $1/2 \ 5/2^+$ | **** |
| $N(1700)$      | 100            | 0.15 | $1/2 \ 3/2^-$ | ***  |

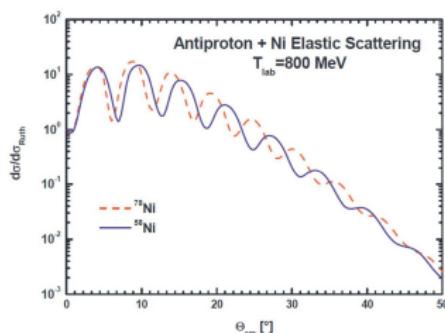
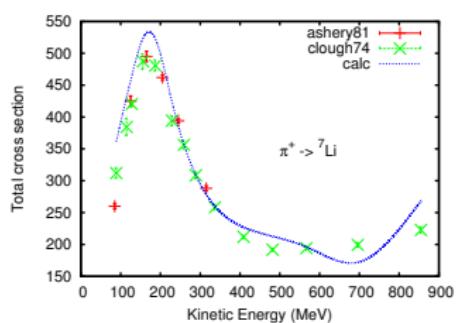
**Table:** Higher  $I = 3/2$  resonance parameters, PDF data.

| Name ( $E_r$ ) | $\Gamma$ [MeV] | R           | $I \ J^P$     | rank |
|----------------|----------------|-------------|---------------|------|
| $\Delta(1600)$ | 350            | 0.10 – 0.25 | $3/2 \ 3/2^+$ | ***  |
| $\Delta(1620)$ | 145            | 0.20 – 0.30 | $3/2 \ 1/2^-$ | **** |
| $\Delta(1700)$ | 300            | 0.10 – 0.20 | $3/2 \ 3/2^-$ | **** |
| $\Delta(1750)$ | 300            | 0.10 – 0.20 | $3/2 \ 1/2^+$ | *    |
| $\Delta(1900)$ | 200            | 0.10 – 0.30 | $3/2 \ 1/2^-$ | **   |
| $\Delta(1905)$ | 330            | 0.09 – 0.15 | $3/2 \ 5/2^+$ | **** |
| $\Delta(1910)$ | 250            | 0.15 – 0.30 | $3/2 \ 1/2^+$ | **** |
| $\Delta(1920)$ | 200            | 0.05 – 0.20 | $3/2 \ 3/2^+$ | ***  |
| $\Delta(1930)$ | 270            | 0.05 – 0.15 | $3/2 \ 5/2^-$ | ***  |
| $\Delta(1940)$ | ~ 200          | 0.05 – 0.15 | $3/2 \ 3/2^-$ | *    |
| $\Delta(1950)$ | 285            | 0.35 – 0.45 | $3/2 \ 7/2^+$ | **** |
| $\Delta(2000)$ | ~ 200          | 0.00 – 0.07 | $3/2 \ 5/2^+$ | **   |

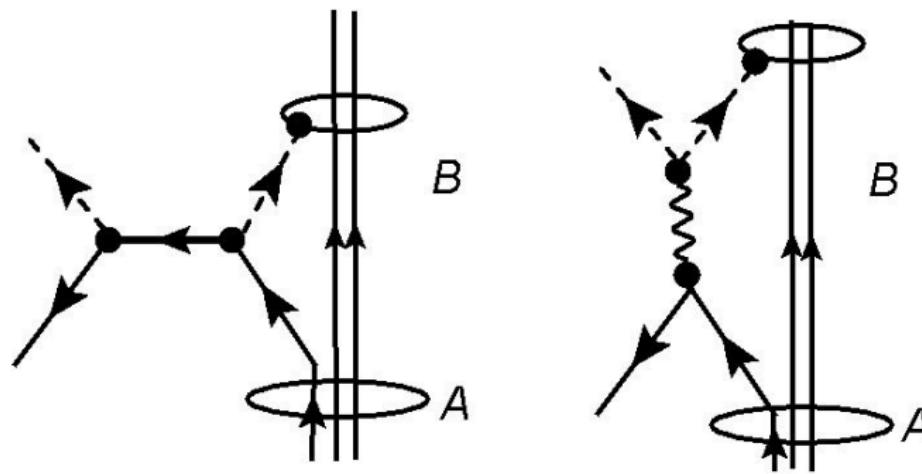




$$M_{\alpha\beta} \approx t_{\bar{N}N \rightarrow 2\pi}(s) \langle \chi_{1\beta}^{(-)} \chi_{2\beta}^{(-)} | \varphi_B | \chi_{\bar{N}A}^{(+)} \rangle \quad (10)$$



# New Feature of Nuclear Annihilation: Single Meson Production



# Summary

- the  $\bar{p}A$  amplitudes are derived in  $t\rho$ -approximation by folding the  $\bar{p}N$  amplitudes with the *HFB*-nucleus densities
- the  $\bar{p}N$  amplitudes are obtained from a (semi-)microscopic model
- hadron production by antiproton annihilation on nuclei
- in progress: meson production as probe for nuclear spectroscopy

In collaboration with J. Haidenbauer (FZ-Juelich).