

Recent Results from CLEO

(selected topics among the most recent results)

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After 29 years CLEO program comes to an end



CESR-b (~10.6 GeV): L=1.2 10^{33} cm⁻²s⁻¹ CESR-c (~4.0 GeV): L=0.7 10^{32} cm⁻²s⁻¹ • CLEO-I (1979-89)

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- CLEO-II (1989-95)
 - Csl calorimeter
- CLEO-II.V (1995-99)
 - Silicon Vertex Detector
- CLEO-III (2000-03)
 - RICH Particle ID
 - New IR & tracking: Silicon, Drift Chamber
- CLEO-c (2003-08)
 - Silicon replaced by ZD inner drift chamber

Last data taken: Run # 234607 Start: 07:38:06 End: 08:00:19 Date: 3/3/2008



CLEO-III/CLEO-c data samples

Detector	Energy	GeV	Lumin	osity	Narrow
	or $Q\overline{Q}$ resonance		fb ⁻¹		resonance
	(targeted	Q q meson)			statistics
CLEO-III	11.227-11.383 Y(5S) (B _s)		0.71		
			0.42		
	Y(4S) (B) + cont	6.2	+ 2.2	
	Y(3S)	+ cont	1.2	+ 0.2	6M
	Y(2S)	+ cont	1.2	+ 0.4	9M
	Y(1S)	+ cont	1.1	+ 0.2	22M
		6.9 - 8.4	0.02		
CLEO-c	4.17 (D _s)		0.61	4	
		3.97-4.26	0.06		
	ψ(3770) (D)		0.818		
	ψ(2S)		0.05	4	27M
		3.673	0.02	1	

Search for light CP-odd Higgs

SM Higgs

 $\widehat{H}_{SM} = \begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix} \implies$

Longitudinal components of W⁺,W⁻,Z⁰ CP-even

 h^0

- Electroweak symmetry breaking
- Direct searches at LEP2: $e^+e^- \rightarrow Z^0 h^0, h^0 \rightarrow b\overline{b} m(h^0) > 114 \text{ GeV}$
- Precision EW measurements: m(h⁰)=89⁺³⁸-28 GeV
- MSSM
 - Fixes hierarchy problem (stabilizes h^0 mass, m(h^0) << M_{Planck} natural)

$$\widehat{H}_{up} = \begin{pmatrix} \varphi_{up}^{+} \\ \varphi_{up}^{0} \end{pmatrix} \quad \widehat{H}_{dn} = \begin{pmatrix} \varphi_{dn}^{0} \\ \varphi_{dn}^{-} \end{pmatrix} \implies \text{L.c. of} \\ W^{+}, W^{-}, Z^{0} \qquad H^{+}, H^{-} \qquad h^{0}, H^{0} \qquad A^{0}$$

- h^0 SM-like. The LEP2 lower limit applies. h^0 , H^0 heavy.
- m(A^0)> m(h^0). A^0 heavy.
- Coupling of $A^{\it 0}$ to down(up-)type fermions $\propto m_f \, tan\beta \,$ ($m_f \, / \, tan\beta$)
- Theoretically expect h^0 to be light: $m(h^0) < 100$ GeV
- μ -problem: $\mu \hat{H}_{up} \hat{H}_{dn}$ why $\mu \ll M_{Planck}$?

 $a_1^0 = A^0 \cos \theta_A - a_c^0 \sin \theta_A$

Search for light CP-odd Higgs

- NMSSM
- Fixes μ -problem: $\lambda \hat{S} \hat{H}_{up} \hat{H}_{dn}$, $\mu_{eff} = \lambda \langle \hat{S} \rangle$ ($\mu << M_{Planck}$ natural) $\hat{H}_{up} = \begin{pmatrix} \varphi_{up}^{\dagger} \\ \varphi_{up}^{0} \end{pmatrix}$ $\hat{H}_{dn} = \begin{pmatrix} \varphi_{dn}^{0} \\ \varphi_{dn}^{\dagger} \end{pmatrix}$ $\hat{S} = (\varphi_{s}^{0}) \Rightarrow$ CP-even CP-odd L.c. of W^{+}, W^{-}, Z^{0} H^{+}, H^{-} h^{0}, H^{0}, h_{s}^{0} A^{0}, a_{s}^{0}
 - Lightest CP-odd Higgs is a mixture:
 - $m(a_1^0) < m(h^0)$ possible
 - If $m(a_1^0) < 2 m_b$ then $a_1^0 \not> b\overline{b}$ and $B(h^0 \rightarrow b\overline{b}) << B(h^0 \rightarrow a_1^0 a_1^0)$ i.e. both $h^0 \& a_1^0$ evade the LEP2 lower mass limit based on bb final state:
 - Relieves "tension" between the direct LEP2 limit and precision EW measurements
 - Relieves "tension" between the direct LEP2 limit and SUSY preference for light h⁰
 - Scenario advocated by Dermisek, Gunion, McElrath [PRD76,051105(R), 2007]
 - Coupling of a_1^0 to down(up)-type fermions
 - $\label{eq:mf} \propto \ m_f \, cos \theta_A tan \beta \ (\ m_f \, cos \theta_A / \, tan \beta \); \quad cos \theta_A \, decreases \ with \ tan \beta$



Search for light CP-odd Higgs



$$\frac{B(\Upsilon(1S) \to \gamma a_1^0)}{B(\Upsilon(1S) \to \mu^+ \mu^-)} = \frac{G_F m_b^2}{\sqrt{2} \pi \alpha} g_d^2 \left\{ 1 - \left(\frac{m(a_1^0)}{m(\Upsilon(1S))}\right)^2 \right\}$$
$$g_d^2 \propto \cos^2 \theta_A \tan^2 \beta$$

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- Good place to produce a₁⁰ with mass m(a₁⁰) < 2 m_b:
 - Favored over lighter mesons by:
 - m_f dependence of the coupling
 - tanβ dependence of the coupling (compared to charmonium)
 - phase-space



Dermisek, Gunion, McElrath [PRD76, 051105(R), 2007]

Search for light CP-odd Higgs

- We have searched for $Y(1S) \rightarrow \gamma a_1^0, a_1^0 \rightarrow \tau^+ \tau^ (a_1^0 \rightarrow \mu^+ \mu^- \text{ for } m(a_1^0) < 2 m_{\tau})$
 - $B(a_1^0 \rightarrow \tau^+ \tau^-)$
 - completely dominates for large tan β (and $m(a_I^{0}) > 2m_{\tau}$)

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Using 21.5M Y(1S) decays

- must compete with $B(a_I^0 \rightarrow c\overline{c})$ for small tan β
- $B(a_1^0 \rightarrow \mu^+ \mu^-)$
 - large below <u>ss</u> threshold (~1 GeV)
 - ~100% for $m(a_1^0)$ slightly above $2m_{\mu}$



3 Σ^+ → $\rho\mu^+\mu^-$ events observed by HyperCP PRL 94, 021801 (2005) He,Tandean, Valencia suggested it could be a a_1^0 PRL 98, 081802 (2007):

Fair amount of fine-tuning to reconcile with the limits from $K \rightarrow \pi \mu^+ \mu^-$

 $g_d \sim O(1)$

→ *observable* $B(Y(1S) \rightarrow \gamma a_1^0)$ [Mangano,Nason Mod.Phys.Lett.A22,1373 (2007)]





• Our $\gamma \tau^+ \tau^-$ limits significantly constrain NMSSM models



 Our γμ⁺μ⁻ limit makes a₁⁰ interpretation of the HyperCP events difficult



- Take V_{cd} = V_{us}=0.2256, determine f_{D,Ds} and compare to lattice QCD calculations:
 - "Calibration point" for lattice QCD calculations needed for B,Bs studies
 - Test of S.M. on f_D/f_{Ds}



- Precision of a few % possible with unquenched calculations
- New predictions of $f_{D^+} = 207 \pm 4 \text{ MeV}$

f_{Ds} = 241±3 MeV by Follana et al HPQCD & UKQCD collaborations (PRL 100, 062002 (2008))



$\mathsf{D^+} \to \mu^+ \nu$ with tagging technique

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- $\bullet \quad e^+e^- \rightarrow \psi(3770) \rightarrow D^+D^-$
 - No fragmentation particles produced
- Reconstruct one D (tag) in several clean hadronic decay modes:
 - Cut on $\Delta E = E_D E_{beam}$
 - Fit $M_{bc} = \sqrt{E_{beam}^2 p_D^2}$ to determine N_{tag}
 - The tag determines momentum of the other D:
 p_{D signal} = **p**_{D tag}
- Find subsample of events with only one additional oppositely charged track within $|\cos\theta| < 0.9$ and no additional photons >250 MeV (to veto D⁺ $\rightarrow \pi^{+}\pi^{0}$)
- Charged track must deposit only minimum ionization in calorimeter [< 300 MeV] (can't use muon system at these low momenta)
- Compute Missing-Mass². If close to zero then almost certainly we have a $\mu^+\nu$ decay.

$$M_{\nu}^{2} = MM^{2} = (E_{beam} - E_{\mu})^{2} - (\vec{p}_{D \text{ signal}} - \vec{p}_{\mu})^{2}$$





 $D^+ \rightarrow \mu^+ \nu$ fit

- $\tau^+\nu/\mu^+\nu$ fixed to SM ratio
 - 149.7±12.0 μv events
 - 28.5 τν
 - $B(D^+ \rightarrow \mu^+ \nu) = (3.86 \pm 0.32 \pm 0.09) \times 10^{-4}$
 - f_D+=(207±8.5±2.5) MeV
- $\tau^+ \nu / \mu^+ \nu$ is allowed to **float**
 - $-153.9\pm13.5 \,\mu\nu$
 - 13.5±15.3 τν
 - $B(D^+ \rightarrow \mu^+ \nu) =$ (3.96±0.35±0.10)x10⁻⁴
 - f_D+=(208.5±9.3±2.5) MeV

CLEO-c Preliminary!





 $D_s^+ \rightarrow \mu^+ \nu$

• Counting tagged D_s^+ events in $e^+e^- \rightarrow D_s^-D_s^*$: 30,848±695

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$$D_s^+ \rightarrow \tau^+ \nu, \tau^+ \rightarrow e^+ \nu \nu$$

316 pb⁻¹ at E_{CM}=4170 MeV

- No $D_s^+ \rightarrow e^+ v$ because of helicity suppression
- MM technique not very useful with 3 neutrinos
- Select events with a tag, electron and then look at extra energy in the calorimeter (includes ~150 MeV γ from $D_s^* \rightarrow \gamma D_s$)
- $B(D_S^+ \rightarrow \tau^+ \nu) =$ (6.17±0.71±0.36)×10⁻²
- f_{Ds}+=(273±16±8) MeV



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 $\Gamma(D_s^+ \rightarrow \tau^+ \nu) / \Gamma(D_s^+ \rightarrow \mu^+ \nu) = 10.3 \pm 1.1$ vs SM(lepton universality)=9.7

Mode	B(%)	f _{Ds} (MeV)
(1) $\mu\nu$ + $\tau\nu$ (fix at SM ratio)	$B^{eff}(D_s \rightarrow \mu \nu) = 0.613 \pm 0.044 \pm 0.020$	268±10±4
(2) μν only	$B(D_s \rightarrow \mu \nu) = 0.600 \pm 0.054 \pm 0.020$	265±12±4
(3) τν, τ→πν	$B(D_s \rightarrow \tau v) = 6.1 \pm 0.9 \pm 0.2$	271±20±4
(4) τν, τ→eνν	$B(D_s \rightarrow \tau v) = 6.17 \pm 0.71 \pm 0.36$	273±16±8
CLEO-c Average of (1) & (4)		269.4±8.2±3.9

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Preliminary !



- We are in close agreement with the Follana et al calculation for f_D+. This gives credence to their calculations.
- The disagreement with f_{Ds} is 3.3σ

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Can we trust lattice QCD calculations?

• If theoretical predictions of f_{Ds}/f_{D} + do not agree with the data, why should we believe f_{Bs}/f_{B} from theory? What does this do to the CKM fits?



If we can trust lattice QCD calculations...

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- New Physics possibilities:
 - Dobrescu, Kronfeld, arXiv:0803.0512



A scalar lepto-quark with electric charge of -1/3

A charged Higgs from Two Higgs doublets model in which one doublet gives masses to u,c quarks and charged leptons, and the other doublet gives masses to down type quarks and top

– Kundu, Nandi, arXiv:0803.1898

Supersymmetry with R-parity violating couplings

M1 Transitions in Charmonium

• Theory:
$$\Gamma(i \xrightarrow{M1} f + \gamma) = \frac{4\alpha e_Q^2}{3m_Q^2}(2J_f + 1)k^3[\mathcal{M}_{if}|]^2$$

$$\mathcal{M}_{if} = \int r^2 dr \, R_{n_i \mathcal{L}_i}(r) j_0(\frac{rk}{2}) R_{n_f \mathcal{L}_f}(r)$$

 $j_0 = 1 - (kr)^2/24 + ..., so in NR limit$

k = 0: $\mathcal{M}_{if} = 1$ $n_i = n_f; L_i = L_f$ "direct" = 0 otherwise "hindered"

$$\Gamma(J/\psi \to \eta_c \gamma) = \frac{16}{3} \alpha e_c^2 \frac{k_\gamma^3}{M_{J/\psi}^2} (1 + \kappa_c) [1 + o(v^2)] \quad \text{(direct)}$$

LQCD 2.0±0.1±0.4 keV Dudek, Edwards, Richards PRD,73,074507(2007) Not much different than the naïve calculation

• Experiment:

 $\begin{array}{l} B(J/\psi \rightarrow \eta_c \gamma) = (1.3 \pm 0.4)\% \text{ Gaiser } et al. \text{ (Crystal Ball), PRD 34, 711 (1986)} \\ \Gamma(J/\psi \rightarrow \eta_c \gamma) = B(J/\psi \rightarrow \eta_c \gamma) \times \Gamma_{tot} = (1.2 \pm 0.3) \text{ keV} \\ \text{Significantly smaller than theoretically predicted!} \end{array}$

CLEO-c method

• $B(J/\psi \rightarrow \gamma \eta_c)$ difficult to measure from inclusive photon spectrum since η_c is broad (25 MeV) and the photon is relatively soft (large background of unknown shape)

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CLEO arXiv:0805.0252 [hep-ex]

 CLEO-c method: DD Threshold 3700 $\eta_{c}' s_{0}$ $B(J/\psi \rightarrow \gamma \eta_c) =$ Mass (MeV) 3300 (MeV) 3300 A / B В B R 'B(J/ψ→γη_c) × B(η_c→X) $B(\Psi' \rightarrow \gamma \eta_c) \times$ $B(\Psi' \rightarrow \gamma \eta_c) \times B(\eta_c \rightarrow X)$ 3100 $\eta_{c} S_{0}$ Measure from 2900 inclusive photon Use exclusive reconstruction of 0-+ 0.1.2++ spectrum. a large number of possible X – reconstructed The photon is final states X (suppresses the exclusive hard, thus backgrounds). decay mode backgrounds are Take the ratio to cancel under control. unknown B($\eta_c \rightarrow X$) anything



Recent CLEO Results M	eson08 T. Skwarnicki 27				
CLEO-c M1 transitions to η_c					
Branching Ratio	CLEO-c	PDG 2006			
$B(\psi(2S) \to \gamma \eta_c)$	$(4.32 \pm 0.16 \pm 0.60) \times 10^{-3}$	$(2.6\pm0.4)\times10^{-3}$			
$B(J/\psi\to\gamma\eta_c)$	$(1.98 \pm 0.09 \pm 0.30) \times 10^{-2}$	$(1.3\pm0.4)\times10^{-2}$			

- The hindered M1 rate $B(\psi(2S) \rightarrow \gamma \eta_c)$:
 - Previous measurements low since they neglected the high energy tail in the signal shape.
 - Theoretical predictions difficult because of the suppressed character of the matrix element.
- The direct M1 rate $B(J/\psi \rightarrow \gamma \eta_c)$:
 - Measured via exclusive event reconstruction. Significantly higher than previously measured by Crystal Ball from inclusive photons.
 - Agrees well with the theoretical expectations. Mystery solved.



Recent CLEO Results Meson08 T. Skwarnicki

New improved determination of h_c mass



CLEO-c arXiv:0805.4599 [hep-ex]

 A factor of ~6 larger statistics than in the initial publication



New improved determination of h_c mass

- Averaging inclusive and exclusive results:
 - $\text{CLEO-c: } m(h_c(1^1P_1)) = 3525.28\pm0.19\pm0.12 \text{ MeV}$
 - vs. spin-averaged m($\chi_c(1^3P_J)$) =3525.30±0.11 MeV
- Thus, hyperfine mass splitting of 1P states is 0.02±0.19±0.13 MeV
 - vs. hyperfine splitting of 1S states of 115 MeV
 - consistent with naïve theoretical prejudice that spin-spin interactions are only short-range and, therefore, not present in 1P states
 - surprisingly small given that corrections to the naïve expectation may be several MeV (see J. M. Richard, Proc. 15th Int. Workshop, DIS 2007 (Munich), Ed. G. Grindhammer and K. Sachs, DESY-PROC-2007-01, p. 849.)

Other recent quarkonium measurements

- Measurement of B(χ_{c0,2}→γγ)=(2.4±0.3±0.2)x10⁻⁴, (3.1±0.3±0.2)x10⁻⁴ arXiv:0803.2869 [hep-ex]
- First measurement of B(J/ψ→γγγ)=(1.2±0.3±0.2)x10⁻⁵ arXiv:0806.0671 [hep-ex]
- Measurements of $B(J/\psi \rightarrow \gamma gg)/B(J/\psi \rightarrow ggg) =$ =0.137±0.001±0.016 arXiv:0806.0315 [hep-ex]
- Update on $B(\psi(2S) \rightarrow X J/\psi)$, $X = \pi^+\pi^-, \pi^0\pi^0, \eta, \pi^0, \gamma\gamma$ via χ_{cJ} arXiv:0804.4432 [hep-ex]
- First measurement of $B(Y(2S) \rightarrow \eta Y(1S)) = (2.1 \pm 0.6 \pm 0.5) \times 10^{-5}$ (preliminary)
- Improved measurements of $B(Y(2S) \rightarrow XY(1S))$, $X = \pi^+ \pi^-, \pi^0 \pi^0$ (preliminary)

Summary

- Search for light CP-odd Higgs in radiative decays of Y(1S):
 - No evidence found
 - Limits in $\gamma \tau \tau$ channel 2 orders of magnitude more stringent than previously available.
 - Limit on $\gamma\mu\mu$ for M($\mu\mu$)=214.3 MeV makes the Higgs interpretation of HyperCP events unlikely
 - Both provide new constraints on NMSSM
- Decay constants of f_D,f_{Ds}:
 - Disagreement with HPQCD lattice QCD calculations on f_{Ds}
 - New physics or problems with HPQCD lattice calculations?
- Measurement of M1 transition rates in charmonium:
 - Significantly different than previously determined
 - Now good agreement with the theory
- Precision measurement of h_c mass:
 - Hyperfine splitting in 1P state very small
- Many other results which I did not have time to describe

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