

# Chiral effective potential with Delta degrees of freedom

**Hermann Krebs**  
Universität Bonn

June 9-th, Meson 2008, Krakow

# Outline

- ChPT and low energy QCD
- $\Delta$ -less few nucleon forces within chiral EFT
  - Expectation from inclusion of  $\Delta$  explicitly
- Isospin-conserving nuclear forces with explicit  $\Delta$
- Isospin-violating nuclear forces with explicit  $\Delta$
- Summary & Outlook

# ChPT and low energy QCD

Spontaneous + explicit (by small quark masses) breaking of chiral symmetry in QCD



Existence of light weakly interacting Goldstone bosons



Systematic description of QCD by ChPT in low energy sector

● Free parameters of QCD

$\alpha_s$ , quark masses

● Free parameters of ChPT

$F_\pi, g_{\pi N}, c_i \dots, d_i \dots, C_S, C_T \dots$



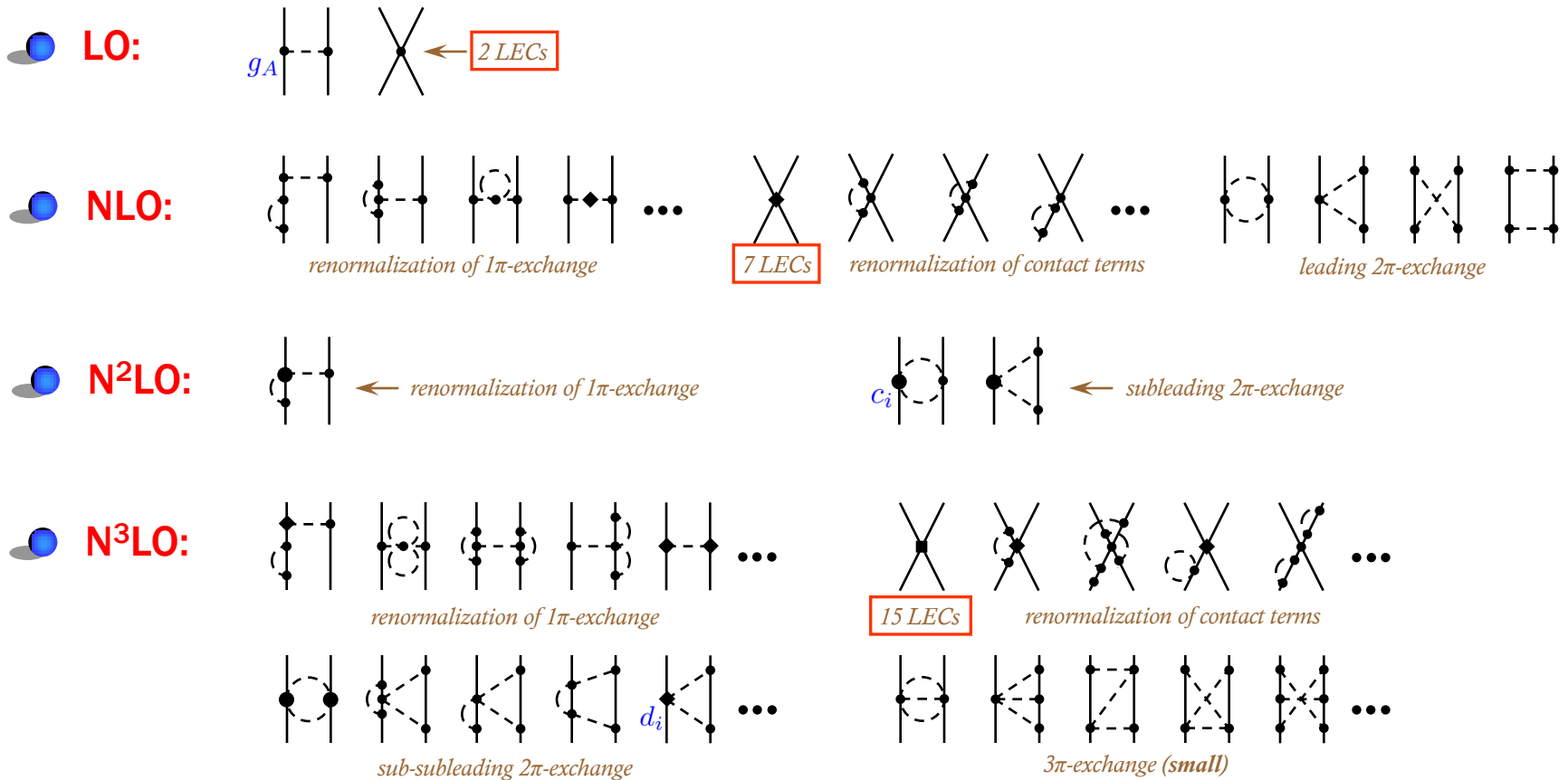
Computable by Lattice QCD / fit to data

# Nucleon-nucleon force up to N<sup>3</sup>LO

Ordóñez et al. '94; Friar & Coon '94; Kaiser et al. '97; Epelbaum et al. '98, '03; Kaiser '99-'01; Higa et al. '03; ...

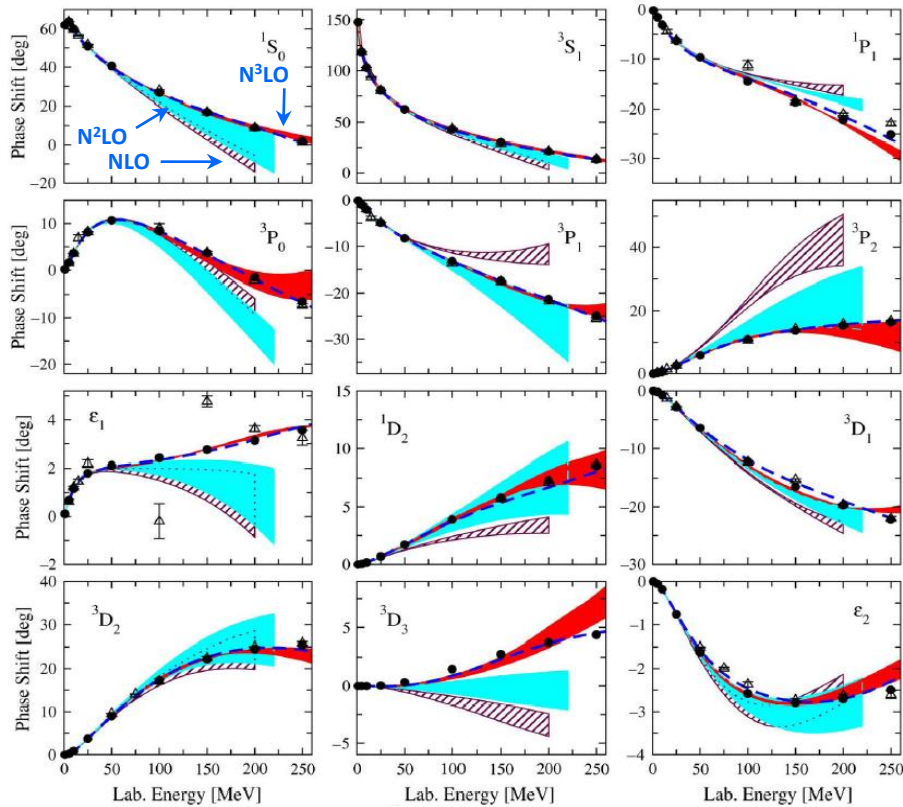
Chiral expansion for the 2N force:

$$V_{2N} = V_{2N}^{(0)} + V_{2N}^{(2)} + V_{2N}^{(3)} + V_{2N}^{(4)} + \dots$$

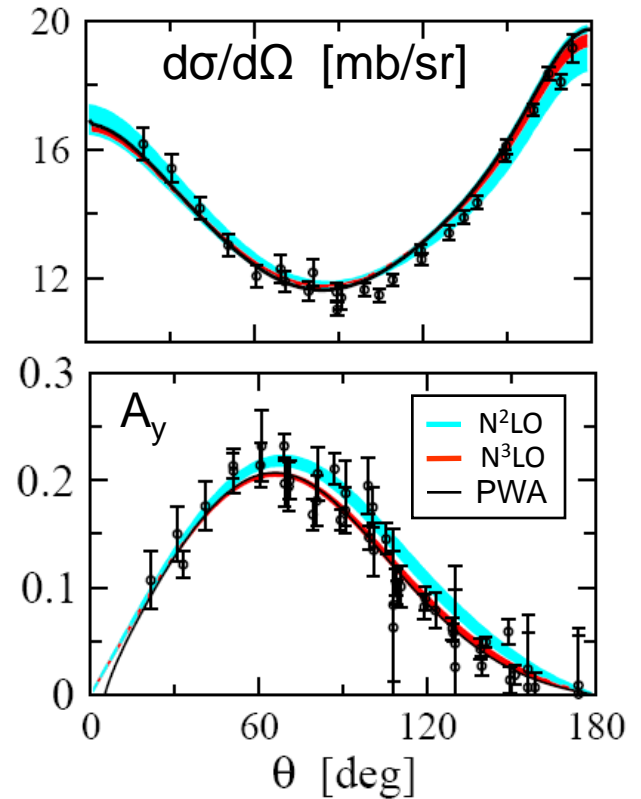


+ 1/m and isospin-breaking corrections...

## Neutron-proton phase shifts up to N<sup>3</sup>LO



## np scattering at 50 MeV



## Deuteron binding energy & asymptotic normalizations $A_S$ and $\eta_d$

	NLO	N <sup>2</sup> LO	N <sup>3</sup> LO	Exp
$E_d$ [MeV]	-2.171 ... -2.186	-2.189 ... -2.202	-2.216 ... -2.223	-2.224575(9)
$A_S$ [ $\text{fm}^{-1/2}$ ]	0.868 ... 0.873	0.874 ... 0.879	0.882 ... 0.883	0.8846(9)
$\eta_d$	0.0256 ... 0.0257	0.0255 ... 0.0256	0.0254 ... 0.0255	0.0256(4)

# Delta-less effective potential

- Standard chiral expansion:  $Q \sim M_\pi \ll \Delta \equiv m_\Delta - m_N = 293 \text{ MeV}$

- Delta contributions encoded in LECs  
(Bernard, Kaiser & Meißner '97)

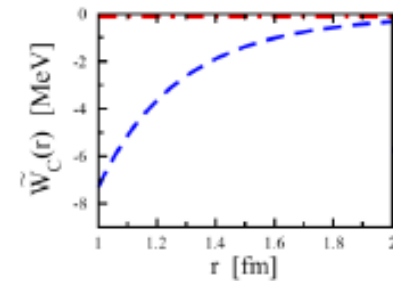
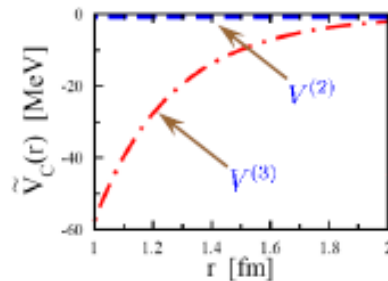
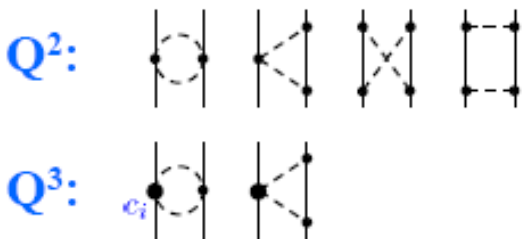


Delta-resonance saturation

$$c_3 = -2c_4 = c_3(\Delta) - \frac{4h_A^2}{9\Delta}$$

Enlargement due to Delta contribution

- Convergence of EFT potential



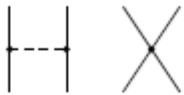



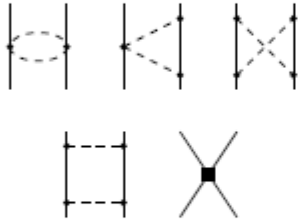
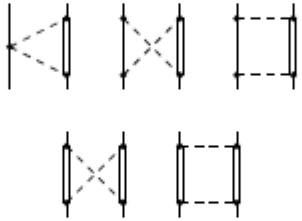



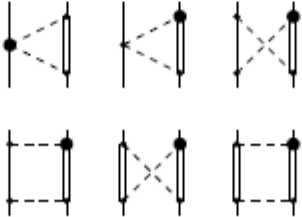
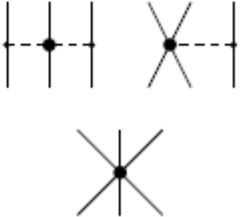

The subleading contribution is bigger than the leading one!

Expectation from inclusion of  $\Delta$  explicitly

- more natural size of LECs
- better convergence
- applicability at higher energies

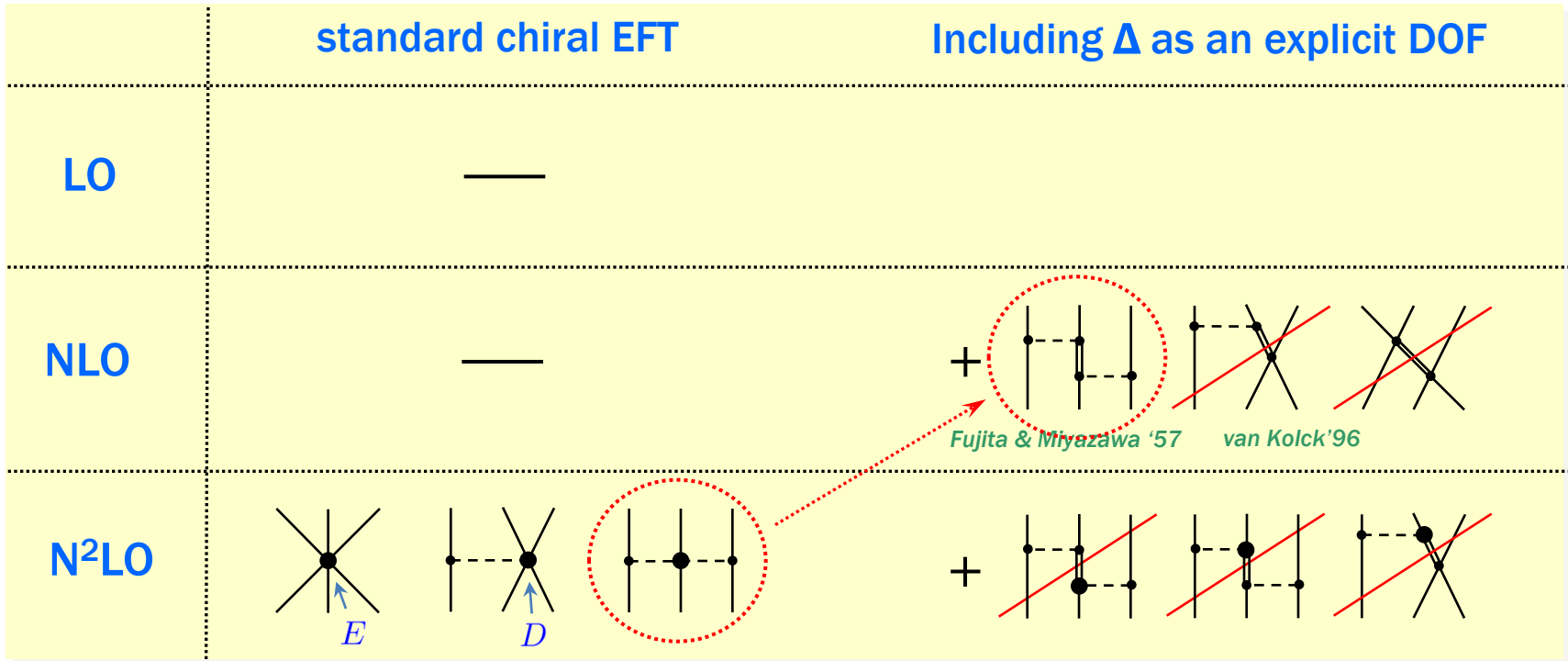
# Few-nucleon forces with the Delta

Isospin-symmetric contributions

	<i>Two-nucleon force</i>		<i>Three-nucleon force</i>	
	$\Delta$ -less EFT	$\Delta$ -contributions	$\Delta$ -less EFT	$\Delta$ -contributions
<i>LO</i>				
<i>NLO</i>		 <i>Ordóñez et al. '96, Kaiser et al. '98</i>		
<i>NNLO</i>		 <i>H.K., Epelbaum &amp; Meißner '07</i>		

# Delta excitations and the three-nucleon force

Epelbaum, H.K., Meißner, *Nucl. Phys. A806 (2008) 65*

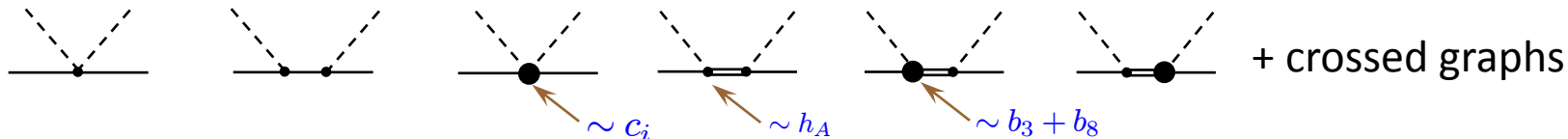


- The LO NNN $\Delta$  contact interaction  $\bar{T}_i^\mu N \bar{N} S_\mu \tau^i N + \text{h.c.}$  vanishes due to the Pauli principle  
 ➡ the LECs  $D$  and  $E$  are not saturated by the delta.
- No contributions from subleading  $2\pi$  –exchange due to  $\partial^0$  at the  $b_3 + b_8$  vertex.
- The entire effect of the  $\Delta$  is given by a partial shift of the N<sup>2</sup>LO TPE 3NF to NLO...



# Fit of LECs to $\pi N$ scattering data

Epelbaum, H.K., Meißner, Eur. Phys. J. A32 (2007) 127



## Two possible values for $h_A$

● fit 1:  $h_A = \frac{3g_A}{2\sqrt{2}} \sim 1.34$  (SU(4), large  $N_c$ )

● fit 2:  $h_A = 1.05$  (Fettes & Meißner '01)

LECs	$Q^2$ , no $\Delta$	$Q^2$ , fit 1	$Q^2$ , fit 2
$c_1$	-0.57	-0.57	-0.57
$c_2$	2.84	-0.25	0.83
$c_3$	-3.87	-0.79	-1.87
$c_4$	2.89	1.33	1.87
$h_A$	-	1.34*	1.05*
$b_3 + b_8$	-	1.40	2.95

## Results of the fit

- Improved description of P-wave parameters when  $\Delta$  is included
- Strongly reduced values for  $c_i$
- Resulting  $c_i$  depend strongly on  $h_A$  while the thresh. param. do not

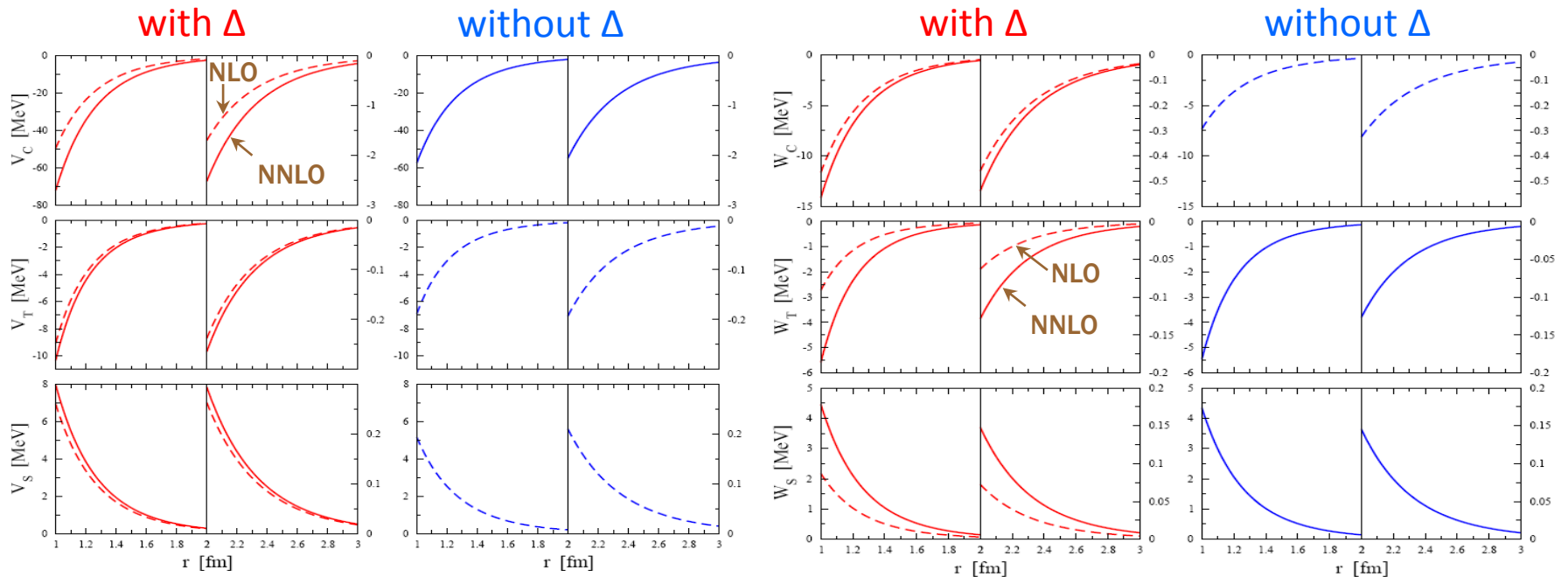
## S- and P-wave threshold parameters

	$Q^2$ , no $\Delta$	$Q^2$ fits 1, 2	EM98
$a_{0+}^+$	0.41	0.41	$0.41 \pm 0.09$
$b_{0+}^+$	-4.46	-4.46	-4.46
$a_{0+}^-$	7.74	7.74	$7.73 \pm 0.06$
$b_{0+}^-$	3.34	3.34	1.56
$a_{1-}^-$	-0.05	-1.32	$-1.19 \pm 0.08$
$a_{1-}^+$	-2.81	-5.30	$-5.46 \pm 0.10$
$a_{1+}^-$	-6.22	-8.45	$-8.22 \pm 0.07$
$a_{1+}^+$	9.68	12.92	$13.13 \pm 0.13$

# NN potential with explicit $\Delta$

$$V_{\text{eff}} = V_C + W_C \vec{\tau}_1 \cdot \vec{\tau}_2 + [V_S + W_S \vec{\tau}_1 \cdot \vec{\tau}_2] \vec{\sigma}_1 \cdot \vec{\sigma}_2 + [V_T + W_T \vec{\tau}_1 \cdot \vec{\tau}_2] (3 \vec{\sigma}_1 \cdot \hat{r} \vec{\sigma}_2 \cdot \hat{r} - \vec{\sigma}_1 \cdot \vec{\sigma}_2)$$

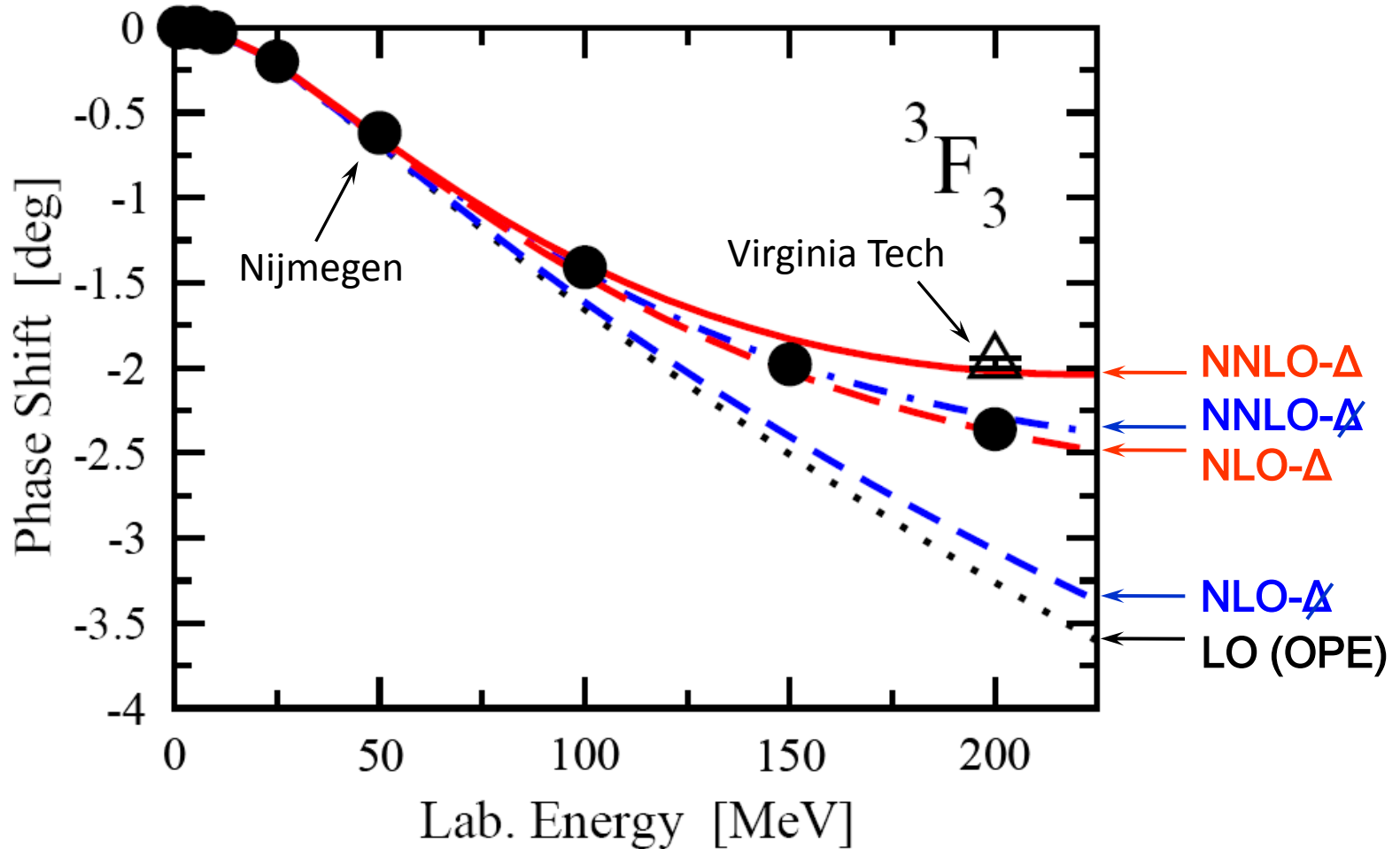
Chiral  $2\pi$ - exchange potential up to NNLO



Advantages when  $\Delta$  is included explicitly

- Dominant contributions already at NLO
- Much better convergence in all potentials

# ${}^3F_3$ partial waves up to NNLO with and without $\Delta$



(calculated in the first Born approximation)

# $\Delta$ -mass splitting in chiral EFT

Epelbaum, H.K., Meißner, Nucl. Phys. A806 (2008) 65

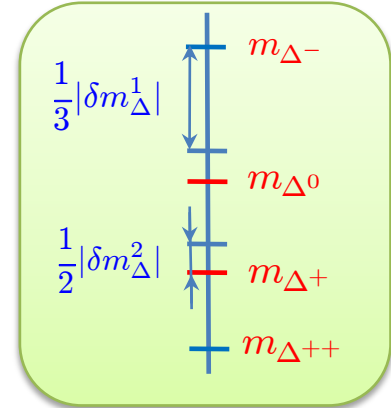
$$\mathcal{L}_{\Delta, \text{mass}}^{\text{LO}} = -\bar{T}_i^\mu \left[ -\delta m_\Delta^1 \frac{1}{2} \tau^3 \delta_{ij} - \delta m_\Delta^2 \frac{3}{4} \delta_{i3} \delta_{j3} \right] g_{\mu\nu} T_j^\nu.$$

Equidistant splitting: strong & em

$$\delta m_\Delta^1 = -4M_\pi^2 \epsilon c_5^\Delta - F_\pi^2 e^2 f_2^\Delta$$

Non-equidistant splitting: em

$$\delta m_\Delta^2 = -\frac{4}{3} F_\pi^2 e^2 f_2^\Delta$$



- Most recent data from the PDG:  $m_{\Delta^{++}} = 1230.80 \pm 0.30$  MeV,  $m_{\Delta^0} = 1233.45 \pm 0.35$  MeV

In addition, PDG's recommended value for the average mass:

$$m_\Delta = \frac{1}{4} (m_{\Delta^{++}} + m_{\Delta^+} + m_{\Delta^0} + m_{\Delta^-}) = \tilde{m}_\Delta + \frac{1}{4} \delta m_\Delta^2 = 1231 \dots 1233 \text{ MeV}$$

On the other hand:  $m_\Delta = 1233.4 \pm 0.4$  MeV (Arndt et al. '06)  $\Rightarrow$  use:  $m_\Delta = 1233$  MeV

$$\Rightarrow \tilde{m}_\Delta = 1233.4 \pm 0.7 \text{ MeV}, \quad \delta m_\Delta^1 = -5.3 \pm 2.0 \text{ MeV}, \quad \delta m_\Delta^2 = -1.7 \pm 2.7 \text{ MeV}$$

- Alternatively, use  $m_{\Delta^{++}}/m_{\Delta^0}$  & the QM relation:  $m_{\Delta^+} - m_{\Delta^0} = m_p - m_n$  (Rubinstein et al. '67)

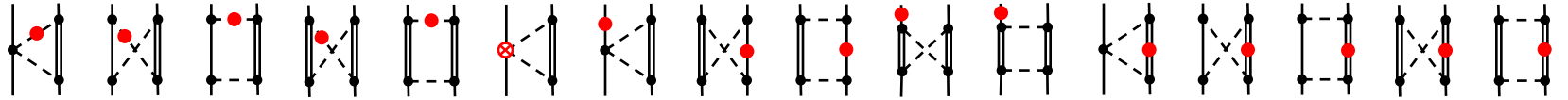
$$\Rightarrow \tilde{m}_\Delta = 1232.7 \pm 0.3 \text{ MeV}, \quad \delta m_\Delta^1 = -3.9 \text{ MeV}, \quad \delta m_\Delta^2 = 0.3 \pm 0.3 \text{ MeV}$$

# Isospin-breaking NN potential

Epelbaum, H.K., Meißner, Phys. Rev. C77 (2008) 034006

2π – exchange contributions with explicit Δ

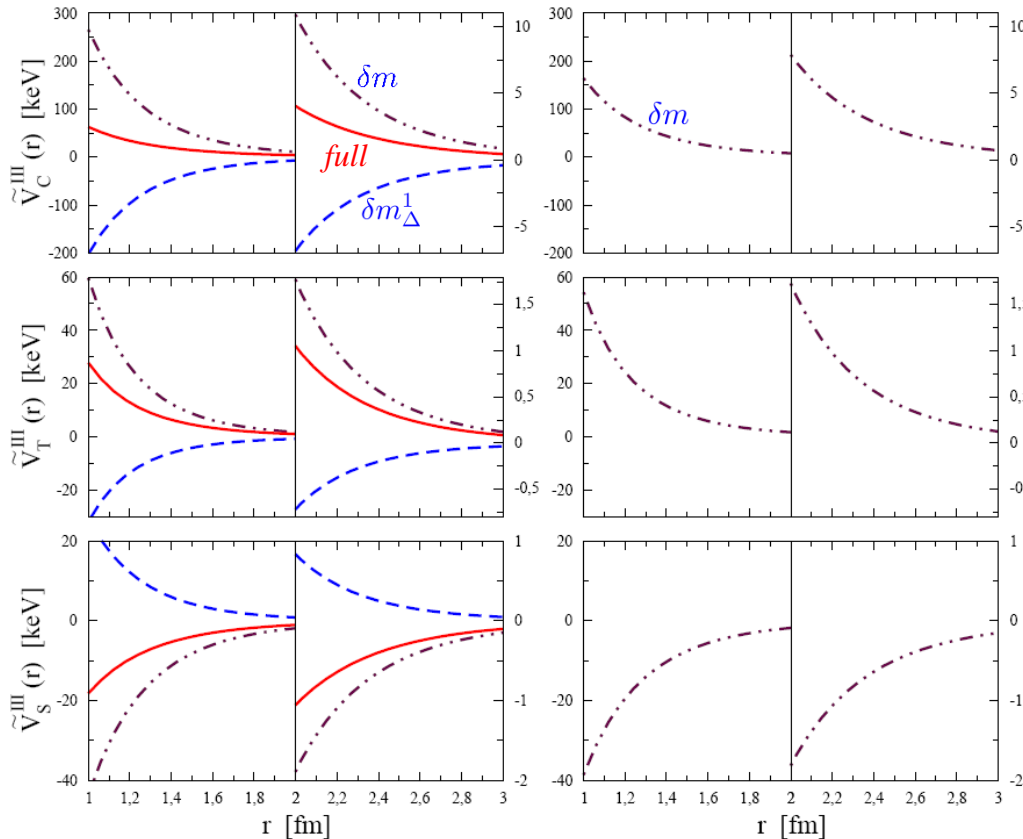
$$V = (\tau_1^3 + \tau_2^3) [V_C^{\text{III}} + V_S^{\text{III}} \vec{\sigma}_1 \cdot \vec{\sigma}_2 + V_T^{\text{III}} \vec{\sigma}_1 \cdot \vec{q} \vec{\sigma}_2 \cdot \vec{q}] + \dots$$



Charge-symmetry-breaking 2π–exchange potential

Δ-full EFT, LO

Δ-less EFT, NLO



- Similar  $\sim \delta m$  contr. to  $\tilde{V}_{S,T}^{\text{III}}$  in the Δ-less and Δ-full EFT
- Sizeable deviation in  $\sim \delta m$  contr. for  $\tilde{V}_C^{\text{III}}$
- Strong cancellations between  $\sim \delta m$  and  $\sim \delta m_{\Delta}^1$  terms



Big contributions beyond the subleading corrections in the Δ-less EFT

# Summary

- Nuclear forces with explicit  $\Delta$  analyzed up to NNLO in chiral EFT
- The unknown LECs fixed from  $\pi$ N S- & P-waves;  $c_i$  strongly reduced
- A more natural convergence pattern for nuclear forces.
  - The most important effects of the  $\Delta$  are well approximated by the saturation of LECs  $c_i$  in the  $\Delta$ -less theory
- Leading isospin-violating  $\Delta$ -contr. to nuclear forces are worked out
  - Delta mass splitting analyzed at LO
  - Found effects beyond resonance saturation of  $c_i$

# Outlook

- Few nucleon observables within the  $\Delta$ -full chiral EFT at NNLO
- Extension to N<sup>3</sup>LO: important effects expected for 3NF and 4NF

